

# DQMC Simulation of the Hubbard Model

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# Overview

## 1. The Physics

*quadratic Hamiltonians, Hubbard model, Hubbard-Stratonovich transformation*

## 2. The Simulation

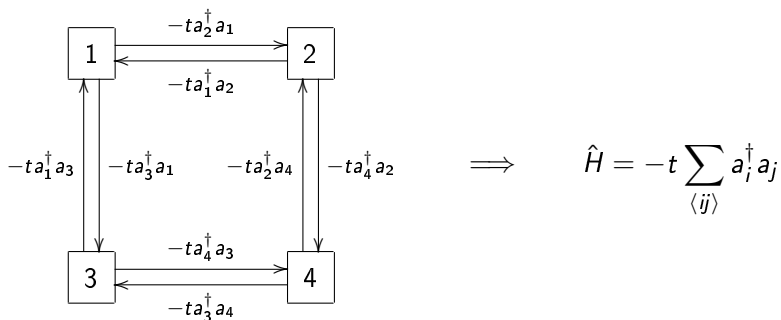
*Markov chain Monte Carlo*

## 3. The System

*Python, NumPy, cProfile, weave, EC2, StarCluster*

# Hopping Hamiltonian

- ▶ Lattice of single-electron states, hopping terms connecting them.



# Hopping Hamiltonian – Quadratic in Operators

- ▶ Adding chemical potential & spin multiplicity,

$$\hat{H} = -t \sum_{\langle ij \rangle} a_{i\sigma}^\dagger a_{j\sigma} - \mu_{i\sigma} \sum_i a_{i\sigma}^\dagger a_{i\sigma}$$
$$= \begin{pmatrix} a_{1\sigma}^\dagger & a_{2\sigma}^\dagger & \dots \end{pmatrix} \begin{pmatrix} h_{11} & h_{12} & \dots \\ h_{21} & h_{22} & \dots \\ \vdots & \vdots & \ddots \end{pmatrix} \begin{pmatrix} a_{1\sigma} \\ a_{2\sigma} \\ \vdots \end{pmatrix}.$$

- ▶ Quadratic  $\Rightarrow$  Exactly solvable!

$$Z = \text{tr} \left[ e^{-\beta \hat{H}} \right] = \det \left[ I + e^{-\beta h} \right].$$

Other rules for Green's functions  $\langle a_{i\sigma}^\dagger a_{j\sigma} \rangle = \frac{\text{tr} [ a_{i\sigma}^\dagger a_{j\sigma} e^{-\beta \hat{H}} ]}{\text{tr} [ e^{-\beta \hat{H}} ]}.$

## Adding Interactions

- ▶ Not too surprising that system is solvable ← no interactions.
- ▶ Simplest possible interaction is on-site repulsion (fixed cost to doubly-occupied site):

$$\hat{H}_{int} = U \sum_i n_{i\uparrow} n_{i\downarrow}$$

...and here's the problem...

$$= U \sum_i a_{i\uparrow}^\dagger a_{i\uparrow} a_{i\downarrow}^\dagger a_{i\downarrow}.$$

- ▶ No longer have a quadratic Hamiltonian.

## Adding Interactions – Discrete H-S Transformation

- ▶ For convenience:

$$\hat{H}_{int} : \quad U \sum_i n_{i\uparrow} n_{i\downarrow} \quad \rightarrow \quad U \sum_i \left( n_{i\uparrow} - \frac{1}{2} \right) \left( n_{i\downarrow} - \frac{1}{2} \right)$$

- ▶ Here's a trick:

$$e^{-\beta U (n_{i\uparrow} - \frac{1}{2})(n_{i\downarrow} - \frac{1}{2})} = C \sum_{s_i} e^{-\beta(\lambda s_i n_{i\uparrow} - \lambda s_i n_{i\downarrow})}$$

$$e^{-\beta \hat{H}_{int}} = C^N \sum_{\{s_i\}} e^{-\beta \sum_i (\lambda s_i n_{i\uparrow} - \lambda s_i n_{i\downarrow})}$$

$$Z = \text{tr} \left[ e^{-\beta \hat{H}_{int}} \right] = C^N \sum_{\{s_i\}} \text{tr} \left[ e^{-\beta \sum_i (\lambda s_i n_{i\uparrow} - \lambda s_i n_{i\downarrow})} \right]$$

- ▶ Sum **quadratic-Hamiltonian** Z's over values of external field!

## Adding Interactions – How does H-S really work?

- ▶  $\uparrow$  electrons and  $\downarrow$  electrons are only connected through “randomly fluctuating” field... how does this create a repulsive interaction?
- ▶ Hubbard-Stratonovich transformation is closely related to mean-field methods.

## Adding Interactions – Trotter Product

- ▶ “Not so fast. . .”

$$e^{-\beta(\hat{H}_0 + \hat{H}_{int})} \neq e^{-\beta\hat{H}_0} e^{-\beta\hat{H}_{int}}.$$

- ▶ Happy resolution – the Trotter product formula!

$$e^{-\beta(\hat{H}_0 + \hat{H}_{int})} = \lim_{\Delta\tau \rightarrow 0} \left( e^{-\Delta\tau\hat{H}_0} e^{-\Delta\tau\hat{H}_{int}} \right)^{\beta/\Delta\tau}.$$

(Our trace-taking techniques survive the breakup.)

- ▶ H-S variables  $s_i$  are now  $s_{i,\tau}$ ; a field in space( $i$ )time.



## Adding Interactions – Making Measurements

- ▶ Summing over all possible  $\{s_{i,\tau}\}$  can give any measurement. For instance, the partition function:

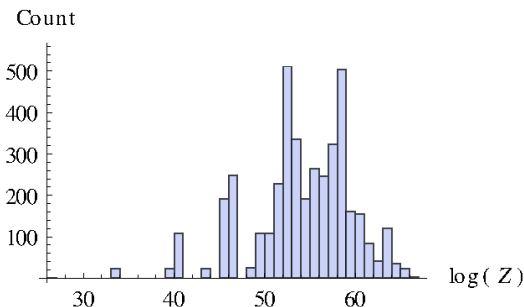
$$\begin{aligned}
 Z &\approx \text{tr} \left[ \left( e^{-\Delta\tau \hat{H}_0} e^{-\Delta\tau \hat{H}_{int}} \right)^{\beta/\Delta\tau} \right] \\
 &= \sum_{\{s_{i,l}\}} \text{tr} \left[ e^{-\Delta\tau \hat{H}_0} e^{-\Delta\tau \hat{H}_{HS}(\{s_{i,L}\})} \dots e^{-\Delta\tau \hat{H}_0} e^{-\Delta\tau \hat{H}_{HS}(\{s_{i,1}\})} \right] \\
 &= \sum_{\{s_{i,l}\}} \det \left[ \underbrace{I + e^{-\Delta\tau h_0} e^{-\Delta\tau h_{HS}(\{s_{i,L}\})} \dots e^{-\Delta\tau h_0} e^{-\Delta\tau h_{HS}(\{s_{i,1}\})}}_{M(\{s_{i,l}\})} \right].
 \end{aligned}$$

- ▶ Or correlation functions:

$$\begin{aligned}
 \left\langle a_{i\sigma}^\dagger a_{j\sigma} \right\rangle_{\{s_{i,l}\}} &= (I - M(\{s_{i,l}\}))^{-1}{}_{ij} \\
 \left\langle a_{i\sigma}^\dagger a_{j\sigma} \right\rangle &= \sum_{\{s_{i,l}\}} \underbrace{\frac{\det [M(\{s_{i,l}\})]}{Z}}_{p(\{s_{i,l}\})} (I - M(\{s_{i,l}\}))^{-1}{}_{ij}.
 \end{aligned}$$

## Markov Chain Monte Carlo – Why?

- ▶ But  $\sum_{\{s_{i,l}\}}$  is only a conceptual possibility
  - ▶  $6 \times 6$  grid, with a *single time-step*? 68,719,476,736 terms
  - ▶ (Each term involves a large matrix  $\times$ ,  $(\cdot)^{-1}$ ,  $\det [\cdot]$ )
- ▶ Random sampling?
  - ▶ Spend too much time in unlikely places:

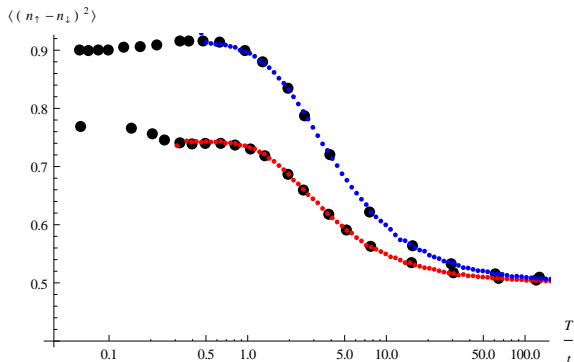


- ▶ (And term computation is still slow)
- ▶ Solution: Markov chain Monte Carlo

# Markov Chain Monte Carlo

- ▶ General idea: Construct a Markov chain with a desired probability distribution as its stationary distribution.
  - ▶ Our case –  $\{s_{i,j}\}$  will take a random walk, by flipping spins with appropriately chosen rates
  - ▶ It will then sample the full distribution accurately
  - ▶ Only requires knowing ratios of probabilities: very easy!
- ▶ Efficient execution comes from keeping track of Green's-function matrices during random walk
  - ▶ test potential flips VERY quickly
  - ▶ update Green's functions in response to a spin-flip using the Sherman-Morrison formula
  - ▶ “wrap” Green's functions from each ( $i$ ) time-step to the next

# Results

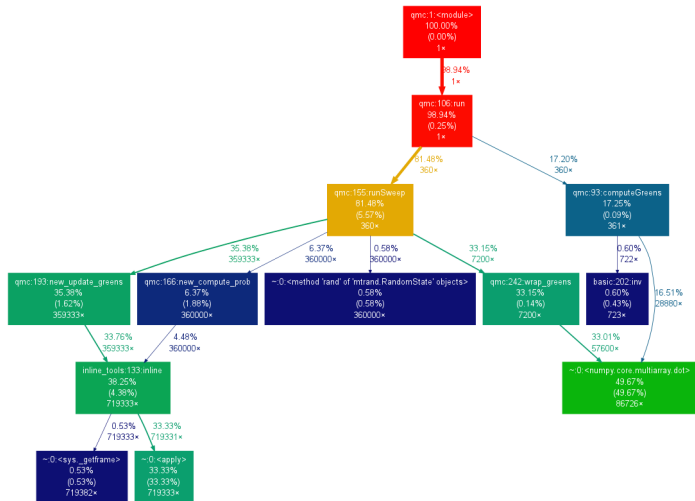


Shown: My results for  $U = 4$  and  $U = 8$ ; vs. Scalettar  
( $6 \times 6$  grid at half-filling)

# The Python Environment

- ▶ **Python:** Modern language with expressive syntax, object-oriented/functional features, and ample libraries.
- ▶ **NumPy:** Adds fast numerical arrays to Python.
- ▶ **SciPy:** Large library of math algorithms based on NumPy.
  
- ▶ Why? I like Python. Wanted to see if it was up to the task.

# Speeding Up Code: cProfile + gprof2dot, line\_profiler



“Premature optimization is the root of all evil.” – Donald Knuth

## Speeding Up Code: weave

- ▶ Put C++ code in Python:

```
def update_G(self, G, new_G, i, Delta_coeff, det_ratio):
    coeff = float(Delta_coeff/det_ratio)
    N = self.N
    code = """
        for (int r = 0; r < N; r++) {
            for (int c = 0; c < N; c++) {
                new_G(r,c) = G(r,c) + G(r, i)*(G(i, c) - (c==i?1:0))*coeff;
            }
        }
    """
    weave.inline(code, ['G', 'N', 'i', 'coeff', 'new_G'],
                type_converters=weave.converters.blitz,
                compiler='gcc')
```

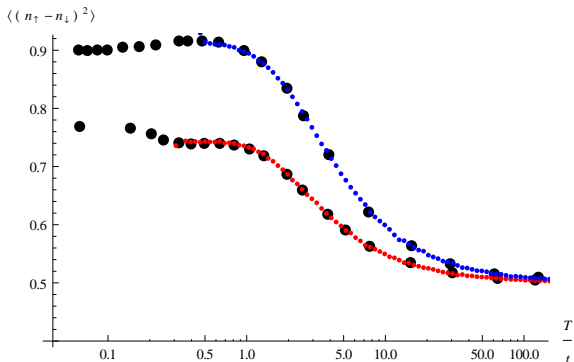
- ▶ Speed up bottlenecks without losing clarity/usability/development-speed of Python.

## Commodity Clusters: Amazon EC2 & StarCluster

- ▶ Amazon EC2 offers virtual machines.
    - ▶ \$0.08/core-hour, on-demand
  - ▶ StarCluster configures everything.
    - ▶ shared filesystem (NFS), queueing system (SGE), etc.
  - ▶ Almost certainly not competitive with academic resources<sup>1</sup>, but a fascinating resource.
    - ▶ Note: SDSC charges \$0.15/core-hour, with a \$2500 minimum.
  - ▶ In 5min computation, with 64 cores...
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# Results



Shown: My results for  $U = 4$  and  $U = 8$ ; vs. Scalettar  
( $6 \times 6$  grid at half-filling)

# Problems & Future Directions

- ▶ Low temperature! Instabilities arise everywhere.
  - ▶ Dealt with so far: adaptive recomputation of Green's functions
  - ▶ Many more to go (fancy multiplication)
- ▶ Scaling
  - ▶ Can the system handle larger lattices? More complex models?
- ▶ Physics
  - ▶ Extend system to explore the universe.